

Ab Initio Calculation of ^2H and ^4He Binding Energies

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ABSTRACT

The binding energies of all hydrogen isotopes have been calculated successfully for the first time in a previous paper [J Fusion Energy, 30 (2011) 377], using only the electric and magnetic Coulomb's laws, without using the hypothetical shell model of the nucleus and its mysterious strong force. In this paper, an elementary calculation gives the order of magnitude of the nuclear interaction. The binding energies of the deuteron and the alpha particle are then calculated by taking into account the proton induced electric dipole in the neutron. The large binding energy per nucleon of ^4He , as compared to that of ^2H , has been explained by a larger electric attraction combined with a lower magnetic repulsion. The binding energies have been calculated without fitting, using only fundamental laws and constants, proving that the nuclear interaction is only electromagnetic.

Keywords: Electromagnetic Moments; Nuclear Forces; Binding Energy Nucleon; Nucleon Interaction; Deuteron; Alpha Particle

1. Introduction

It is known since one century that radium releases a huge energy, one million times larger than any combustion energy, according to Pierre Curie. Let us compare the separation energy ratio of a proton from a neutron over that of an electron from a proton. The measured value of this ratio is $2.2 \text{ MeV}/13.6 \text{ eV} = 160,000$, less than the million expected because the deuteron is lightly bound. The radius ratio of the hydrogen atom ($0.53 \times 10^{-10} \text{ m}$) over the proton's ($\approx 1 \text{ fm} = 10^{-15} \text{ m}$) is above 50,000. As far as I know, no theoretical formula for the proton radius exists. A simple approach using the proton Compton radius ($0.210 \times 10^{-15} \text{ m}$) as the proton radius, leads to a formula giving an order of magnitude of the nuclear binding energy, αmc^2 , predicting the nuclear energy to be around 1% of the mass energy. Therefore, the nuclear to chemical energy ratio is shown to be 250,000, not far from the experimental value 160,000 of the ratio between the binding energies of the deuteron and the hydrogen atom. More precise calculations using the electromagnetic neutron-proton interaction confirm this rough approximation as it will be shown below.

2. Simple Approach to the Nuclear Interaction

The Bohr radius of the hydrogen atom (Figure 1) is:

$$a_0 = \frac{\hbar}{\alpha m_e c} = 0.53 \times 10^{-10} \text{ m}. \quad (1)$$

where

$$\alpha = \frac{e^2}{2\epsilon_0 hc} = \frac{1}{137} \quad (2)$$

is the fine structure constant, h , Planck's constant, m_e , the electron mass and c , the light velocity. No theoretical formula existing for the radius of a nucleon [1], we shall use the proton Compton radius R_p instead although it is four times smaller than the experimentally evaluated value of the proton radius:

$$R_p = \frac{\hbar}{m_p c} = 0.21 \times 10^{-15} \text{ m} \quad (3)$$

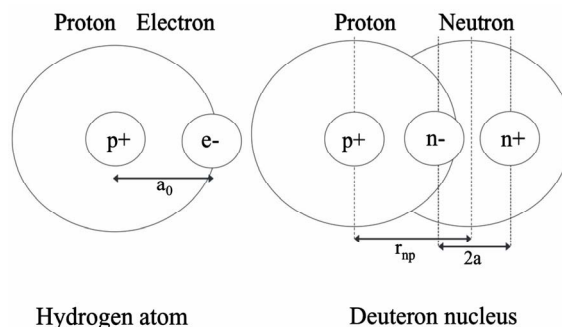


Figure 1. Comparison between the electron and neutron distances from a proton. $-a_0$ is the Bohr radius, r_p is the proton-neutron separation distance and $2a$ the separation distance between the electric charges of the neutron's electric dipole induced by the proton.

The ratio of the Bohr radius a_0 (1) over the proton Compton radius R_p (3) is:

$$\frac{a_0}{R_p} = \frac{m_p}{m_e \alpha} = 1836 \times 137 = 250,000 \quad (4)$$

where m_p is the proton mass. The Bohr formula of the binding energy of the fundamental state of the hydrogen atom is:

$$\frac{1}{2} \alpha^2 m_e c^2 = 13.6 \text{ eV} \quad (5)$$

Newton's law of gravitation and Coulomb's law of electricity are the only forces of nature having a potential energy inversely proportional to the distance. Assuming that it is the same for the nuclear interaction, the ratio a_0/R_p from Equation (4) is the ratio of nuclear and chemical energies. Multiplying it by the hydrogen atom binding energy (5) we obtain the total binding energy of the deuteron:

$$\frac{1}{2} \alpha m_p c^2 = 3.5 \text{ MeV} \quad (6)$$

This value is larger than the experimental binding energy of the deuteron, 2.2 MeV. The binding energies per nucleon varying from 1 MeV for the deuteron to 9 MeV for iron, we may say that the order of magnitude of the nuclear binding energy per nucleon is around

$$\alpha m_p c^2 = 1/137 \times 939 = 6.85 \text{ MeV} \quad (7)$$

This value is, coincidentally, almost that of ^4He , 7.05 MeV. This simple calculation based on the hypothesis of an inverse distance law for the nuclear potential predicts, as it is well known, the nuclear energy to be

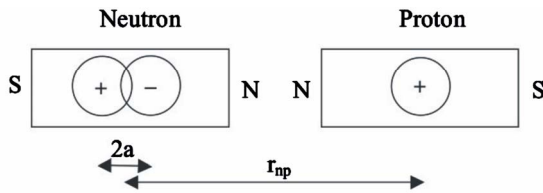


Figure 2. Schematic deuteron structure. r_{np} is the neutron-proton separation distance. $2a$ is the electric dipole moment separation distance. The electric dipoles are at the centers of the nucleons. At large internucleon distances the electric neutron-proton interaction energy potential is that of a dipole, in $1/r^2$, explaining the apparent neutrality of the neutron at large distances from a proton. At a short internucleon distance, the positive charge of the neutron is repulsed by the proton and its negative charge is attracted by the potential in $1/r$ as in any electrostatic induction [2]. The magnetic moments of the neutron and the proton are opposite North-North (or South-South) and collinear (not antiparallel as usually assumed), thus producing a repulsive interaction potential in $1/r^3$ equilibrating the electrostatic attractive potential in $1/r$.

around $1\% \approx \alpha = \frac{1}{137}$ of the mass energy.

3. Electromagnetic Interaction in a Nucleus

In contrast with the Bohr planetary model of the atom, the nucleus has no nucleus and thus no fixed axis of rotation, the center of mass of the nucleus being not precisely defined. It is usually admitted that the centrifugal force is equilibrated by the mysterious strong force. It is assumed here that a static equilibrium between attractive electrostatic and repulsive magnetic forces exists.

The usual dipole and polarizability formulas being invalid in a non-uniform electric field e.g. between a neutron and a nearby proton, the original Coulomb's law for point charges is used. The electrostatic interaction in the nucleus is due to the opposite elementary electric charges separated in a neutron by a nearby proton, inducing an electric dipole. The magnetic interaction between the proton and the neutron is due, in the deuteron, to the collinear and opposite magnetic moments of the nucleons (Figure 2).

In the ^4He nucleus (Figure 3), the electromagnetic interaction works with the same principle as for ^2H with two differences. First, the electrical dipole in a neutron is induced by two protons, implying a larger elec-

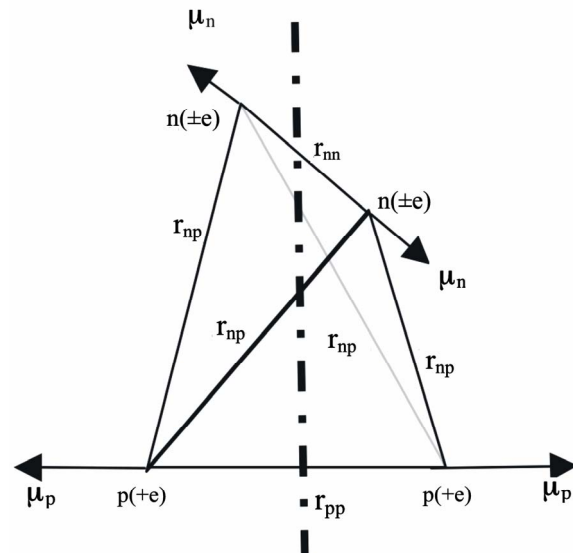


Figure 3. Tetrahedral alpha particle. The proton contains one elementary charge and the neutron two equal and opposite charges. All electric charges are assumed to be equal, in absolute value, to the elementary charge e . All magnetic moments are horizontal and perpendicular to the symmetry axis. The magnetic moments of the protons being opposite and collinear, the resulting magnetic moment is zero. Same thing for the neutrons. The magnetic moments of the protons are perpendicular to those of the neutrons. The tetrahedron is assumed to be regular: $r_{nn} = r_{pp} = r_{np}$.

trical interaction energy than in ${}^2\text{H}$. Second, the magnetic dipoles are inclined at 60° with respect to the vertices, implying a lower neutron-proton repulsion. This explains the high binding energy of the α particle. Only universal and fundamental constants are used: elementary electrical charge e , neutron and proton magnetic moments μ_n , μ_p , vacuum electric permittivity ϵ_0 , magnetic permeability μ_0 , light speed c or, equivalently, fine structure constant α , proton mass m_p , neutron and proton Landé factors g_n , g_p , proton Compton radius R_p .

3.1. Electric Charges in the Neutron

If the neutron had no charge, its electrostatic energy would be zero and it should be lighter than the proton. This is wrong: the neutron is heavier than the proton [2] by 1.29 MeV. The mass of the electron is 0.51 MeV. The kinetic energy of the electron, the proton and the electron antineutrino is the Q-value, difference between the masses before and after the free neutron β decay:

$$Q = 1.29 - 0.51 = 0.78 \text{ MeV} \quad (8)$$

Gamow [3] suggested the electron-proton model where the neutron contains two opposite elementary electric point charges $+e$ and $-e$ and the proton only one positive point charge. The presence of electric charges in the neutron is known since the discovery of its magnetic moment [4].

3.2. Electric Dipoles

When a proton approaches a neutron, the positive electric charge of the neutron is repulsed by the proton while the negative charge is attracted, creating a dipole. This dipole is not permanent, it disappears when the proton goes far away from the neutron.

Let us see first the potential energy of a permanent dipole (only **three collinear charges** are considered here, $+e$ for the proton, $-e$ and $+e$ for the neutron):

$$U_e^{2H} = \frac{e^2}{4\pi\epsilon_0} \left(\frac{1}{r_{np} + a} - \frac{1}{r_{np} - a} \right) \approx -\frac{e^2}{4\pi\epsilon_0} \frac{2a}{r_{np}^2} \quad (9)$$

r_{np} is the distance between the proton and the dipole center. $2a$ is the separation distance between the induced charges. The approximate dipole formula, at the right, is valid only when $a \ll r_{np}$, in a quasi-uniform electric field. It will not be used here, where the separation distance between the neutron and the proton is comparable to the separation distance between the electric charges of the neutron. When the proton is bound to the neutron, the proton induced electric dipole, combined with the proton electric charge, becomes the quadrupole moment of the deuteron, $Q = 0.288 \text{ fm}^2 = (0.54 \text{ fm})^2$ meaning that the distance between the electric charges is

comparable to the nucleon size. The neutron dipole is not permanent: it is induced by the proton providing the energy needed to create the dipole. Therefore, the energy provided by the proton when it approaches the neutron has to be added to the self-energy of the dipole. Both energies being given by the same formula, the total interaction energy of the neutron and the proton is twice that of a permanent dipole:

$$U_e^{2H} = \frac{e^2}{4\pi\epsilon_0} \left(\frac{2}{r_{np} + a} - \frac{2}{r_{np} - a} \right) \quad (10)$$

An almost equivalent assumption is to assume that the neutron behaves like an isolated neutral conductor and that the proton is a point charge near to the neutron:

“When you bring a positive charge up to a conducting sphere, the positive charge attracts negative charges to the side closer to itself and leaves positive charges on the surface of the far side” [2].

This phenomenon, investigated by Faraday who called it “Electrification by Induction” [5], should also happen in the “not so neutral neutron” [6] even if its conductivity and charges are unknown. We have partial induction but we may consider it, in a first approximation, as a total induction. The charges induced by $+e$ (the proton) will thus be the same elementary charges $-e$ and $+e$ as above.

Formula (10) may be written differently (conversion formulas between (10) and (11) are given in the appendix):

$$U_e^{2H} = \alpha m_p c^2 \left(\frac{2R_p}{r_{np} + a} - \frac{2R_p}{r_{np} - a} \right) \quad (11)$$

3.3. Magnetic Dipoles

In the deuteron ${}^2\text{H}$, the neutron and the proton have opposite and, by reason of symmetry, **collinear magnetic moments**, resulting in a repulsive force. The magnetic moment of ${}^4\text{He}$ being zero, the protons (same thing for the neutrons) are paired, collinear and opposite. Assuming provisionally a regular tetrahedron, the magnetic moments of the protons and the neutrons are, also by reason of symmetry, perpendicular and inclined at 60° with respect to the vertices of the assumed regular tetrahedron (**Figure 3**):

$$U_m^{ij} = \frac{\mu_0}{4\pi r_{ij}^3} \left[\boldsymbol{\mu}_i \cdot \boldsymbol{\mu}_j - \frac{3(\boldsymbol{\mu}_i \cdot \mathbf{r}_{ij})(\boldsymbol{\mu}_j \cdot \mathbf{r}_{ij})}{r_{ij}^2} \right] \quad (12)$$

4. The Coulomb's Laws

The electric Coulomb's potential energy:

$$U_e = \frac{e^2}{4\pi\epsilon_0 r} \quad (13)$$

may be written equivalently for nuclear physics (see appendix):

$$U_e = \alpha m_p c^2 \frac{R_p}{r} \quad (14)$$

where α is the fine structure constant, m_p the proton mass, c the light velocity and R_p the proton Compton radius. A similar expression exists for the magnetic Coulomb's potential as we shall see below.

4.1. The Electromagnetic Potential Energy in a Nucleus

The sum of the electrostatic interaction energy potential between electric charges e_i and e_j separated by $r_{ij} \pm a_{ij}$, plus the magnetic interaction energy potential between nucleons with magnetic moments μ_i and μ_j , separated by r_{ij} is [5,7,8]:

$$U_{em} = \sum_i \sum_{i \neq j} \frac{\epsilon e_i e_j}{4\pi\epsilon_0 (r_{ij} \pm a_{ij})} + \sum_i \sum_{i \neq j} \frac{\mu_0}{4\pi r_{ij}^3} \left[\mu_i \cdot \mu_j - \frac{3(\mu_i \cdot r_{ij})(\mu_j \cdot r_{ij})}{r_{ij}^2} \right] \quad (15)$$

where $r_{ij} \pm a_{ij}$ is the separation distance between the electric charges (always equal to e in absolute value in this paper) and r_{ij} the separation vector between the nucleons magnetic moments. We have $\epsilon = 2$ if the interaction is between a neutron and a proton because of the induced dipole, needing twice the energy of a permanent dipole; otherwise it is 1 between protons or 0 between neutrons. This general Formula (15) shows that the electric Coulomb potential is attractive or repulsive depending on the sign of the product of the interacting electric charges. The magnetic potential energy is attractive or repulsive depending on the relative orientation and position of the magnetic moments of the nucleons. Using the fundamental constants shown in the appendix, the general potential energy Formula (15) may be converted into:

$$U_e = \alpha m_p c^2 \sum_i \sum_{i \neq j} \left[\frac{\epsilon e_i e_j R_p}{e^2 (r_{ij} \pm a_{ij})} \right] \quad (16)$$

for the electrostatic potential and

$$U_m = \alpha m_p c^2 \sum_i \sum_{i \neq j} \left[\frac{|g_i g_j|}{16} \left(\frac{R_p}{r_{ij}} \right)^3 S_{ij} \right] \quad (17)$$

for the magnetic potential where the g 's are the Landé factors. S_{ij} , positive for a magnetic repulsion and negative for a magnetic attraction, is the tensor operator [9],

r_{ij} is the internucleon vector and μ_i , μ_j are the interacting magnetic moments of the nucleons i and j :

$$S_{ij} = \cos(\mu_i, \mu_j) - 3 \cos(\mu_i, r_{ij}) \cos(\mu_j, r_{ij}) \quad (18)$$

The electromagnetic nuclear potential is the product of $\alpha m_p c^2 = 6.8 \text{ MeV}$ and a purely numerical function to be determined. The total electromagnetic potential is:

$$U_{em} = U_e + U_m \quad (19)$$

4.2. Deuteron Electromagnetic Energy Potential

The deuteron has one proton (one positive charge) and one neutron (two equal and opposite charges) (Figure 2) with three electric interactions and one magnetic interaction between the proton and the neutron. The proton interacts with the induced $+e$ and $-e$ charges of the neutron dipole. As seen above, the energy of the neutron electric dipole has to be added because it is not pre-existent, thus multiplying by 2 the proton-neutron electrostatic interaction. Indeed the exact formula of the dipole potential is the same as for the interaction between a point charge and two opposite charges. The tensor operator is, for collinear and opposite magnetic moments, according to Formula (18):

$$S_{np} = \cos(\mu_n, \mu_p) - 3 \cos(\mu_n, r_{np}) \cos(\mu_p, r_{np}) = -1 - 3 \times 1 \times (-1) = 2 \quad (20)$$

Formula (19) becomes:

$$U_{em}^{2H} = \alpha m_p c^2 \left[\frac{2R_p}{r_{np} + a} - \frac{2R_p}{r_{np} - a} + 2 \frac{|g_n g_p|}{16} \left(\frac{R_p}{r_{np}} \right)^3 \right] \quad (21)$$

Numerically, in MeV, where the neutron-proton separation distance vector $|r_{np}|$ and the electric dipole moment separation distance $2a$ (Figures 1 and 2), are in fm:

$$U_{em}^{2H} = 6.85 \left(\frac{2 \times 0.210}{r_{np} + a} - \frac{2 \times 0.210}{r_{np} - a} \right) + 6.85 \times 2 \times \frac{3.83 \times 5.59}{16} \left(\frac{0.210}{r_{np}} \right)^3 \quad (22)$$

The minimum of the potential giving the binding energy for one bond, it has to be divided by two to obtain the binding energy per nucleon of the deuteron.

4.3. α Particle Electromagnetic Energy Potential

The helium ${}^4\text{He}$ (Figure 3) has one nn , one pp and 4 np bonds. The nn bond electrostatic energy may be neglected because there is probably no electric interaction between the neutrons. The magnetic moments of the

protons being collinear and opposite along the same edge, there is electric and magnetic repulsion between the protons. Because there are two protons inducing each neutron, the electrostatic potential is multiplied by 2 with respect to that of the deuteron, doubling the electrostatic attraction. Therefore, the coefficient of R_p of the electrostatic terms of the neutron-proton interaction is 4 (instead of 2 for the deuteron where there is induction by only one proton on one neutron). The coefficient $1/4$ is due to the single proton-proton bond for 4 neutron-proton bonds. The electric potential per nucleon between a neutron and a proton plus between protons (between neutrons it should be 0) is, from Equation (16):

$$U_e^{4\text{He}} = \alpha m_p c^2 \left(\frac{4R_p}{r_{np} + a} - \frac{4R_p}{r_{np} - a} + \frac{1}{4} \times \frac{R_p}{r_{pp}} \right) \quad (23)$$

We have also to take into account the inclination between the magnetic moments. The magnetic moments of the proton and the neutron being perpendicular, the first term of (16) is zero. Being inclined at 60° with respect to their r_{np} bond, their cosines are $1/2$. The projections on r_{np} are opposite and make an angle of 120° with r_{np} . Therefore, Formula (18) becomes:

$$S_{np} = \cos(\boldsymbol{\mu}_n, \boldsymbol{\mu}_p) - 3 \cos(\boldsymbol{\mu}_n, \mathbf{r}_{np}) \cos(\boldsymbol{\mu}_p, \mathbf{r}_{np}) \quad (24)$$

$$= 0 - 3 \times \frac{1}{2} \times \left(-\frac{1}{2} \right) = \frac{3}{4}$$

The general Formula (18) gives thus a factor $S_{np} = 3/4$ instead of 2 for the deuteron (21). The np magnetic component of ${}^4\text{He}$ is thus $3/8$ times smaller than in the deuteron. The magnetic moments of the protons being parallel and perpendicular to the straight line joining them, we have

$$S_{pp} = S_{nn} = \cos(\boldsymbol{\mu}_p, \boldsymbol{\mu}_p) - 3 \cos(\boldsymbol{\mu}_p, \mathbf{r}_{pp}) \cos(\boldsymbol{\mu}_p, \mathbf{r}_{pp}) \quad (25)$$

$$= 1 - 3 \times 0 = 1$$

According to Formulas (24) and (25), the magnetic component of the electromagnetic potential of ${}^4\text{He}$, for one bond (or one nucleon) is thus:

$$U_m^{4\text{He}} = \alpha m_p c^2 \left(\frac{3}{4} \times \frac{|g_n g_p|}{16} + \frac{1}{4} \times \frac{g_p^2 + g_n^2}{16} \right) \left(\frac{R_p}{r_{np}} \right)^3 \quad (26)$$

The electromagnetic potential, for one bond of the α particle, is:

$$U_{em}^{4\text{He}} = U_e^{4\text{He}} + U_m^{4\text{He}}$$

$$= \alpha m_p c^2 \left(\frac{4R_p}{r_{np} + a} - \frac{4R_p}{r_{np} - a} + \frac{1}{4} \times \frac{R_p}{r_{pp}} \right) \quad (27)$$

$$+ \alpha m_p c^2 \left(\frac{3}{4} \times \frac{|g_n g_p|}{16} + \frac{1}{4} \times \frac{g_p^2 + g_n^2}{16} \right) \left(\frac{R_p}{r_{np}} \right)^3$$

Numerically,

$$U_{em}^{4\text{He}} = 6.85 \left(\frac{4 \times 0.210}{r_{np} + a} - \frac{4 \times 0.210}{r_{np} - a} \right) \quad (28)$$

$$+ 6.85 \times \frac{1}{4} \times \frac{3 \times 3.83 \times 5.59 + 5.59^2 + 3.83^2}{16} \left(\frac{0.210}{r_{np}} \right)^3$$

The minimum of the potential is the binding energy per nucleon, the number of neutron-proton bonds being equal to the number of nucleons.

5. Binding Energies of ${}^2\text{H}$ and ${}^4\text{He}$

Because of the two variables, r_{np} , the neutron-proton separation distance and $2a$ the dipole moment separation distance, the binding energy cannot be derived analytically as was done in an earlier paper [10]. It has been solved graphically by trial and error until finding the energy potential minimum for both the internucleon distance and the neutron electric dipole. **Figure 4** shows that there is in fact no real minimum, only an inflection point contrarily to the first calculation of the deuteron where the minimum was clear [10]. This may be amended by taking into account the finite structure of the electric charges, perhaps by using numerical or electric image methods. The potential energies of ${}^2\text{H}$ and ${}^4\text{He}$ corresponding to Formulas (22) and (28) are on **Figure 4**, showing the agreement between experiment and theory.

The binding energy per nucleon of ${}^2\text{H}$ is found to be 1.05 MeV practically the experimental value 1.1 MeV.

The calculated binding energy of ${}^4\text{He}$ is 6.2 MeV significantly lower than the experimental value, 7.1 MeV.

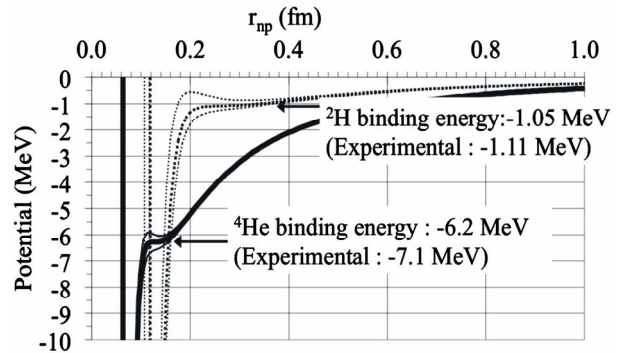


Figure 4. ${}^2\text{H}$ and ${}^4\text{He}$ electromagnetic energy potentials per nucleon from Formulas (22) and (28). The binding energies are obtained graphically by trial and error from the electric and magnetic Coulomb's potentials by varying the neutron proton distance r_{np} and the dipole separation $2a$ until a relative minimum is obtained, unfortunately a flat spot, due to the Coulomb singularity. A true minimum was obtained for the deuteron when the positive charge of the neutron is neglected [10]. Only fundamental constants are used, no fitting or *ad hoc* parameters.

This discrepancy may be solved with a tetrahedron less symmetrical. Between protons and neutrons there is electric and magnetic repulsion. There is probably no electric interaction between neutrons. Its magnetic moment is smaller than that of the proton. Thus, the electromagnetic repulsion between neutrons being smaller than between protons, the separation distance should be also smaller between neutrons than between protons.

The minimum of the potential (**Figure 4**) occurs at $r_{np} \approx 0.25$ fm for ${}^2\text{H}$ and at $r_{np} \approx 0.15$ fm for ${}^4\text{He}$. The empirical nuclear potentials of the literature give larger values, around 0.5 fm.

6. Discussion

The potential has a minimum when the positive charge of the neutron is neglected as was shown earlier [10]. When the dipole is taken into account, the potential has no real minimum, only a flat spot. This inflection point is due to the Coulomb singularity when the distance a between the electric charges approaches the separation distance r between the centers of the nucleons.

The fine structure constant α appears when the electron charge e and the absolute dielectric permittivity of classical vacuum ϵ_0 are replaced by the proton mass, the light velocity c and the proton Compton radius $R_p = 0.21$ fm. Although R_p is 4 times smaller than the measured proton radius, 0.87 fm and half the usual value of the radial minimum of the nucleon-nucleon potential energy, usually around $r_{np} = 0.5$ fm. The calculations give good results for the binding energies without needing any *ad hoc* parameter.

In the deuteron, the magnetic moments of the proton and the neutron are opposite because the magnetic moment of the deuteron is, approximately, the difference between the absolute values of the proton and neutron magnetic moments. When the magnetic moments are collinear, the proton and the neutron rotate around their common axis, stabilized by the gyroscopic effect due to the nucleon spin.

The calculated binding energy of ${}^4\text{He}$ is still 20% too weak. Indeed, the symmetry of ${}^4\text{He}$ is not that of a regular tetrahedron because there are two kinds of nucleons with different electric and magnetic properties. A more precise calculation will be performed as soon as possible, taking into account a lower symmetry of the tetrahedron.

7. Results

The following results have been obtained by applying the electromagnetic theory to the atomic nucleus:

- The nuclear attraction between a neutron and a proton is the electrostatic induction of a proton on a nearby neutron.
- The soft core is the repulsion between the magnetic moments of the nucleons.
- The calculated binding energies of the deuteron and of the α particle agree satisfactorily with the experimental data.
- A nuclear equivalent of the Rydberg constant, $\frac{1}{2}\alpha^2 m_e c^2$ has been found to be $\alpha m_p c^2$.
- The ratio between nuclear and chemical energy is discovered to be $\frac{m_p}{m_e \alpha}$.

8. Conclusion

The electric and magnetic Coulomb's laws applied to the nucleons (without orbital angular momenta) suffice to predict quantitatively the nuclear interaction as the achievement of the ${}^2\text{H}$ and ${}^4\text{He}$ binding energies proves it, the calculations being easily verifiable. The agreement with experiment confirms the electromagnetic nature of the nuclear interaction found in a preceding paper [10]. Taking into account the neglected interactions in ${}^4\text{He}$ should enhance its binding energy precision. In contrast, the hypotheses of charge independence, strong force and shell model are unable to calculate ${}^2\text{H}$, the simplest nucleus beyond the proton and *a fortiori* ${}^4\text{He}$. It is hoped to generalize the electromagnetic approach to all nuclei.

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Appendix: Fundamental Constants Used

- Fine structure constant

$$\alpha = \frac{e^2}{2\epsilon_0 hc} = \frac{1}{137} \quad (29)$$

- Proton Compton radius

$$R_p = \frac{\hbar}{m_p c} = 0.21 \text{ fm} \quad (30)$$

- Nuclear magneton

$$\mu_N = \frac{e\hbar}{2m_p} = \frac{ecR_p}{2} \text{ J.T}^{-1} \quad (31)$$

- Magnetic moments of the neutron and the proton μ_n and μ_p and their corresponding Landé factors, $g_n = -3.826$ and $g_p = 5.585$, are related by

$$\mu_i = \frac{g_i}{2} \mu_N = g_i R_p \frac{ec}{4} \text{ J.T}^{-1} \quad (32)$$

where i means n or p .

- Relation between vacuum dielectric permittivity and magnetic permeability

$$\epsilon_0 \mu_0 c^2 = 1 \quad (33)$$

- Fundamental constants of the nuclear energy potential
Electrostatic attraction:

$$\alpha m_p c^2 = \frac{938}{137} = \frac{e^2}{4\pi\epsilon_0 R_p} = 6.85 \text{ MeV} \quad (34)$$

4% weaker than the α particle binding energy per nucleon (-7.1 MeV).

Magnetic repulsion:

$$\frac{\mu_0 |\mu_n \mu_p|}{4\pi R_p^3} = \alpha m_p c^2 \frac{|g_n g_p|}{16} = 9.15 \text{ MeV} \quad (35)$$

where $g_n = -3.826$ and $g_p = 5.585$.